



FIG. 11. The function $(W - W_I)T^{-1}$ versus T for Ge as a means of determining the extrapolated lattice thermal resistivity at high temperatures.

Thus

$$K = K_\theta + K_e + K_r, \tag{7}$$

where K_θ is the phonon, K_e is the electronic, and K_r is the radiative or photon contribution to K .

Let us first consider the lattice thermal conductivity K_θ . At high temperatures we have to consider the relaxation times τ_I and τ_U given in Eqs. (3) and (4), but we can neglect τ_B in Eq. (6). Equation (4) gives the relaxation time for 3-phonon umklapp processes. However, for $300^\circ\text{K} < T < 1681^\circ\text{K}$, we are in the range of T comparable to or greater than the Debye temperature θ for both Si and Ge. In this range, it may be necessary to consider the relaxation times for four-phonon processes, as Pomeranchuk⁴⁵⁻⁴⁷ has pointed out. He gives a relaxation time for these higher order (H) processes as:

$$\tau_H^{-1} = B_H \omega^2 T^2, \tag{8}$$

with B_H a constant. K_θ can be evaluated for $T > \theta$ from Eqs. (1), (3), (4), and (8) if τ_C^{-1} is taken as

$$\tau_C^{-1} = \tau_U^{-1} + \tau_H^{-1} + \tau_I^{-1}.$$

In the region where $T > \theta$ the quantity x in Eq. (1) is small, and the integral simplifies to

$$K = \frac{k}{2\pi^2 v} \left(\frac{kT}{\hbar} \right)^3 \int_0^{\theta/T} \tau_C x^2 dx. \tag{9}$$

Also the exponential factor in B_U disappears to make B_U temperature-independent. Thus

$$\tau_C^{-1} = (B_U T + B_H T^2) \omega^2 + A \omega^4. \tag{10}$$

For $T \geq \theta$ the isotope scattering is much less important than the phonon-phonon scattering. In this limit Eqs. (9) and (10) can be reduced by the method used by Ambegaokar,⁴⁸ to

$$K_\theta^{-1} \equiv W_\theta = W_U + W_H + W_I,$$

where

$$\begin{aligned} W_U &= \pi v h B_U T / \theta k^2, \\ W_H &= \pi v h B_H T^2 / \theta k^2, \\ W_I &= 4\pi^2 V_0 \theta \Gamma / h v^2. \end{aligned} \tag{11}$$

This reduction requires $W_\theta \gg W_I$. This condition is fulfilled for Si and Ge at high temperatures. The only really unknown quantity in Eq. (11) is B_H . The quantities B_U and B_H can be evaluated experimentally from a plot of $(W_\theta - W_I)T^{-1}$ versus T . The quantity W_I is, except for a difference of a factor of 12 in the definition of Γ , the same as that given by Ambegaokar.⁴⁸ For Si one obtains $W_I = 0.033$ cm deg/W. For Ge the value is $W_I = 0.17$ cm deg/W using $V_0 = 2.26 \times 10^{-23}$ cm³, $\theta = 395^\circ\text{K}$,⁴⁹ $\Gamma = 4.90 \times 10^{-5}$,⁵⁰ and $v = 3.94 \times 10^5$ cm/sec.

⁴⁸ V. Ambegaokar, Phys. Rev. 114, 488 (1959).

⁴⁹ P. Flubacher, A. J. Leadbetter, and J. A. Morrison, Phil. Mag. 4, 273 (1959).

⁵⁰ D. Strominger, J. M. Hollander, and G. T. Seaborg, Rev. Mod. Phys. 30, 585 (1958).

⁴⁵ I. Pomeranchuk, Phys. Rev. 60, 820 (1941).
⁴⁶ I. Pomeranchuk, J. Phys. USSR 4, 259 (1941).
⁴⁷ I. Pomeranchuk, J. Phys. USSR 7, 197 (1943).